## ORIGINS OF TERAHERTZ PHOTORESPONSE IN GRAPHENE TRANSISTORS: THEORY AND EXPERIMENT

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2024

## **Outline:**

- 1. Introduction:
- What is THz radiation?
- THz Detectors. Why graphene (CNT) based
- 2. The main mechanisms of THz radiation detection by graphene-based FET devices
- 3. Asymmetric devices based on graphene:
- Device fabrication and experimental setup
- Device characterization
- Response of asymmetric devices based on graphene on THz radiation
- Contribution of plasmonic response to the detection of sub-terahertz radiation using graphene based devices
- 4. Conclusions



# **THz range is important for:**

- Security
- Medicine
- Astronomy
  - Etc.



Visualization of the hidden objects



THz image of liver

Astronomy



Detectors. Why nanosturtures:

- -Sensitive
- Fast
- Energy efficient
- Etc.

**Detectors. Why graphene (CNT) based:** 

-Gapless graphene has strong interband absorption at all frequencies

- High room-temperature mobility
- -Geometric control of the band structure
  - Easy to fabricate



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$$E(k) = E_0 + 2\gamma_1 \left( \frac{c}{c} \cos(\vec{k}\vec{a}_1) + \cos(\vec{k}\vec{a}_2) + \cos(\vec{k}(\vec{a}_1 - \vec{a}_2)) \right)$$
  
$$\pm \gamma_0 \sqrt{3 + 2\cos(\vec{k}\vec{a}_1) + 2\cos(\vec{k}\vec{a}_2) + 2\cos(\vec{k}(\vec{a}_1 - \vec{a}_2))}$$

## **Detectors.** Why graphene (CNT) based:

Figure 1.8: The low-energy band structure of monolayer graphene (1.57) taking into account nearest-neighbor hopping with parameter  $\gamma_0 = 3.033$  eV, nearest-neighbor overlap parameter  $s_0 = 0.129$ , and orbital energy  $s_0 = 0.129$  [11]. The plot shows the bands calculated in the vicinity of the first Brillouin zone, with conduction and valence bands touching at six corners of the Brillouin zone, two of them are labeled  $K_+$  and  $K_-$ . Label  $\Gamma$ indicates the center of the Brillouin zone. Adopted from[27]

## sorption

## - High room-temperature mobility

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Plasmonics forms a major part of the fascinating field of *nanophotonics*, which explores how electromagnetic fields can be confined over dimensions on the order of or smaller than the wavelength.

Plasmonics: Fundamentals and Applications Authors: Maier, Stefan Alexander Springer, 2007

## **Detectors. Why graphene (CNT) based:**

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#### What do we mean by graphene THz detector

 $\langle I \rangle$ 



$$\begin{split} \left(\delta V \cos \omega t\right) \Big\rangle_{T} \approx \\ \frac{dI}{dV} \delta V \left\langle \cos \omega t \right\rangle_{T} + \frac{1}{2} \frac{d^{2}I}{dV^{2}} \delta V^{2} \left\langle \cos^{2} \omega t \right\rangle_{T} \\ = \frac{1}{4} \frac{d^{2}I}{dV^{2}} \delta V^{2} \end{split}$$

Study of detection mechanisms is the study of nonlinearities



Current due to built-in field at the junction OR rectification due to diode nonlinearity Voltage due to Resistance change temperature gradients due to in nonuniformly doped overall device heating channel Formation of standing plasma waves in the device channel

Figure from: F.H.L. Koppens, T. Mueller, P. Avouris, A.C. Ferrari, M.S. Vitiello, M. Polini, "Photodetectors based on graphene, other two-dimensional materials and hybrid systems" Nature nanotechnology, 9, 780-793 (2014)

## Two types of asymmetric graphene based structures



## Dyakonov – Shur configuration



**Difference** in workfunction results in formation of a p-n junction along the channel. Schottky diode

Asymmetric boundary conditions:  $V_S = V_{\theta} \cos(\omega t); I_D = \theta$ Result in a DC voltage signal as the device is exposed to radiation

#### **Photo-thermoelectric effect in graphene**

In case of a **photothermoelectric effect** an non-uniform doping of the channel and non-uniform heating of the channel results in onset of a DC voltage proportional to increase of the electron temperature



#### Graphene advantages for hot-electron photothermoelectric detection:

o-Gapless graphene has strong interband absorption at all frequencies.
o-The electronic heat capacity of single-layer graphene is much lower than in bulk materials, resulting in a larger change in temperature for the same absorbed energy
o- The photothermoelectric effect has a picosecond response time, set by the electron– phonon relaxation rate

#### Nano Lett. 16, 6988 (2016)

#### Photo-thermoelectric effect in graphene

#### Graphene photothermoelectric detector. Principle of operation





Graphene photothermoelectric detector device fabrication and principle of operation. (a-e) Lithographic sequence used to produce the graphene terahertz detector.(f) Optical micrograph showing electrical contacts and (inset) atomic force micrograph showing bimetallic contacts connected to an exfoliated graphene layer. (g-k) Schematic of the principle components during device operation. (g) Cross-sectional view of the device. (h-j) Profiles across the device of (h) electron temperature T(x), (i) Fermi level EF(x), (j) Seebeck coefficient S(x) and (k) potential gradient

#### \*Nature Nanotechnology 9, 814–819 (2014)



Broadband thermoelectric responsivity of graphene photothermoelectric detector. (a,d) Electrical conductance, (b,e) responsivity to Joule heating, and (c,f) responsivity to radiation as a function of gate voltage for the device shown in Fig. 1f at room temperature and in ambient environment.

#### **Detection, mixing, and frequency multiplication of terahertz radiation by two-dimensional** electronic fluid\*



The basic equations describing the two dimensional electronic fluid the are relationship between the surface carrier concentration and gate voltage swing, the equation of motion, and the continuity equation

$$n_{s} = CU/e$$
$$\frac{\partial v}{\partial t} + v\frac{\partial v}{\partial x} + \frac{e}{m}\frac{\partial U}{\partial x} + \frac{v}{\tau} = 0$$
$$\frac{\partial U}{\partial t} + \frac{\partial (Uv)}{\partial x} = 0.$$

#### The boundary conditions

$$U(0,t) = U_o + U_a \cos \omega t \quad \text{for} \quad x = 0$$
  
$$j(L,t) = 0 \quad \text{for} \quad x = L$$

Here

$$\beta = \frac{2\omega\tau}{\sqrt{1+(\omega\tau)^2}}$$

\*IEEE TRANSACTIONS ON ELECTRON DEVICES, VOL. 43, NO. 3, MARCH 1996

#### Detection, mixing, and frequency multiplication of terahertz radiation by two-dimensional electronic fluid



$$\frac{\Delta U}{U_o} = \frac{1}{4} \left(\frac{U_a}{U_o}\right)^2 f(\omega)$$

 $f(\omega) = 1 + \beta - \frac{1 + \beta \cos\left(2k'_o L\right)}{\sinh^2\left(k''_o L\right) + \cos^2\left(k'_o L\right)}.$ 

Here

$$eta = rac{2\omega au}{\sqrt{1+(\omega au)^2}}$$



\*IEEE TRANSACTIONS ON ELECTRON DEVICES, VOL. 43, NO. 3, MARCH 1996

#### **Graphene field-effect transistors as room-temperature terahertz detectors\***



а

50 µm

1µm

300 nm

b

Due to high roomtemperature mobility (up to **10 000 cm<sup>2</sup>/(V s)** on SiO<sub>2</sub> graphene is promising for THz FET photodetectors.

#### Diffusive transport model

 $\omega \tau_{ee} << 1$ 

 $\omega \tau$  tr<<1

 $j(x,t) = \sigma E(x,t) = -\sigma \frac{\partial V_{G}(x,t)}{\partial x}$ 

in conjunction with the continuity equation:

Vgate

 $\frac{\partial [-en(x,t)]}{\partial t} + \frac{\partial j(x,t)}{\partial x} = 0$ 

The total carrier density -en(x, t) is modulated by the gate voltage  $V_G(x, t)$  according to

 $\Delta u = \frac{U_a^2}{4} \frac{1}{\sigma(U_a)} \left. \frac{\mathrm{d}\sigma(V_G)}{\mathrm{d}V_G} \right|_{V_G = U_0}$ 

$$-en(x,t) = CV_{G}(x,t)$$

С

#### **Graphene field-effect transistors as room-temperature terahertz detectors\***



$$\Delta u = \frac{U_a^2}{4} \frac{1}{\sigma(U_0)} \left. \frac{\mathrm{d}\sigma(V_\mathrm{G})}{\mathrm{d}V_\mathrm{G}} \right|_{V_\mathrm{G}=U_0}$$

The minimum RT **NEP** is 200nWHz<sup>-1/2</sup> for SLG and almost one order of magnitude lower (30nWHz<sup>-1/2</sup>) for BLG.



#### NATURE MATERIALS; VOL 11 ;OCTOBER 2012 j





## Highest mobility devices probed in DS configuration

D. A. Bandurin, I. Gayduchenko, et al., 112, 141101, (2018)







(a) Room temperature the response voltage is a linear function of the radiation power

More importantly, R(Vg) follows the F(Vg) with F = dG/dVg \* R



At 77 K peak responsivity increases as the radiation power goes down R(Vg) still follows the F(Vg)



Temperature evolution of responsivity is different for different gate voltages

## Heating-induced photoresponce



$$U_{PTE} = -\int SdT \approx S(T_{S} - T_{D})$$

 $S \approx -\frac{\pi^2 k_{\rm B}^2 T}{3e} \frac{1}{\sigma} \frac{d\sigma}{dE_{\rm F}}$  is the Seebeck coefficient

$$F = \frac{1}{\sigma} \frac{d\sigma}{dV_{\rm bg}}$$

Heated region





## Heating-induced photoresponce



**Temperature evolution of responsivity is different for in case of PURE PTE accounting for the p-n junction at the Ti-Graphene interface (***simulations performed by D. Svintsov*)

## **Dyakonov-Shur photoresponse**



$$\Delta U_{\rm DS} = \frac{U_{\rm a}^{2}}{4} \frac{1}{\sigma} \frac{d\sigma}{dV_{\rm bg}} g(\omega) \propto U_{\rm a}^{2} F, \quad (3)$$

where  $U_a$  is the antenna voltage, L is the channel length,  $g(\omega) = (\sinh^2 kL - \sin^2 kL)/(\sinh^2 kL + \cos^2 kL)$  is the "form factor" depending on the wave number  $k = \sqrt{\omega/(2s^2\tau)}$  of the overdamped plasma wave characterized by the group velocity  $s = \sqrt{4\alpha_{ee}d\sqrt{\pi n}}$ , n is the carrier density and  $\alpha_{ee} \sim 1$  is the parameter describing the strength of e-e interactions in





## Conclusions

- FET based on graphene encapsulated in hBN can serve as high-responsivity THz detector
- Terahertz detection in graphene FETs is a combination of resistive self-mixing, photothermoelectric effects and p-n junction rectification
- Hydrodynamics, strong electron-electron and electron-hole scattering are good for photodetection