

# Quantum fluctuations in superconducting nanostructures

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# Outline

## 1. Introduction

**Fluctuations in quasi-1D superconductors**  
**Quantum Phase Slip concept**

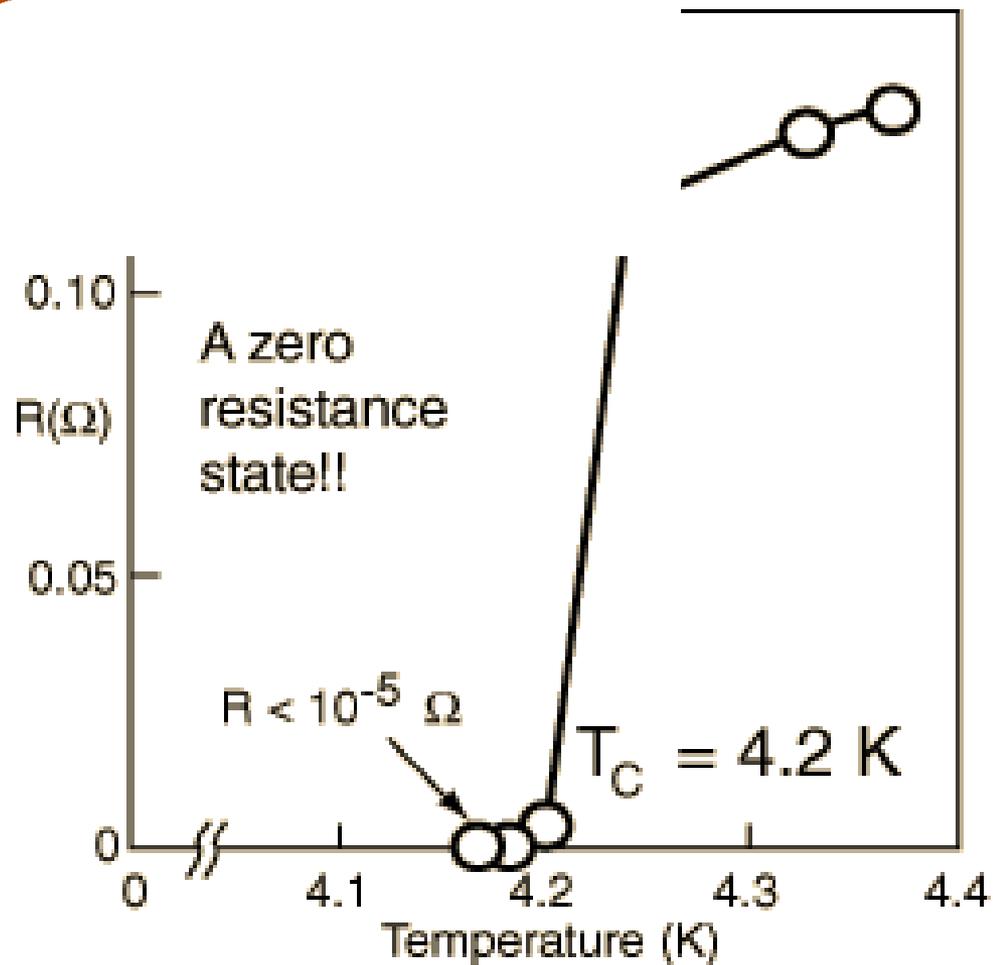
## 2. Applications

**Junctionless Cooper pair transistor**  
**Quantum standard of electric current**

## 3. Conclusions

# RITD transition

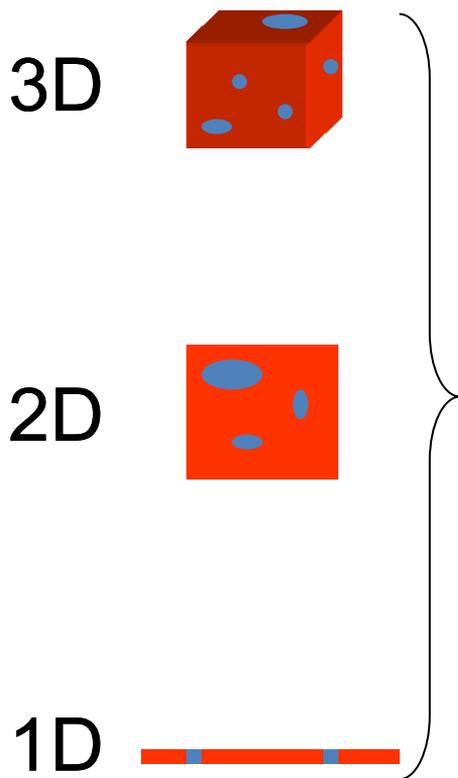
H. K. Onnes,  
Commun. Phys. Lab.12,120, (1911)



# Fluctuations vs. system dimensionality

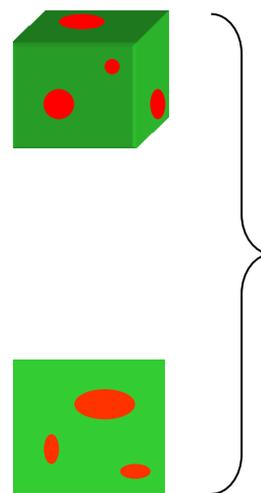
For a superconductor dimensionality is set by the temperature - dependent coherence length  $\xi(T)$ .

**N** normal metal  
**S** superconductor

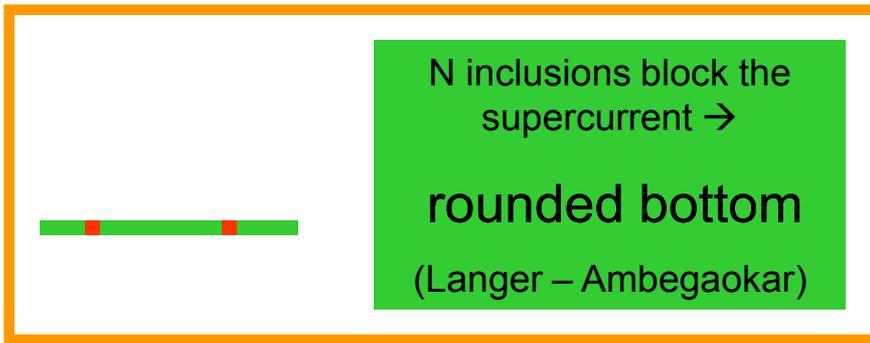


S inclusions reduce the total system resistance  $\rightarrow$   
**rounded top**  
 $\sigma_{\text{FLUCT}} \sim (T - T_c)^{-(2-D/2)}$   
(Aslamazov - Larkin)

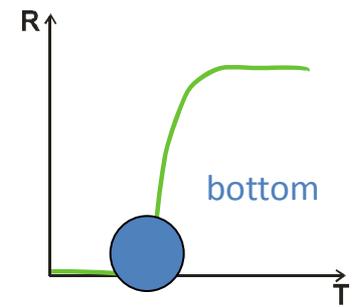
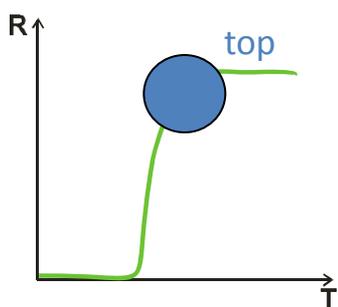
Topic of the talk



no contribution of N inclusions: normal current is shunted by supercurrent  $\rightarrow$   
**abrupt bottom**



N inclusions block the supercurrent  $\rightarrow$   
**rounded bottom**  
(Langer - Ambegaokar)



# Fluctuations in a 1D superconductor

Long 1D wire of cross section  $\sigma$    $\updownarrow \quad v\sigma < \xi(T) \ll L$

If the wire is infinitely long, there is always a finite probability that in some fragment(s) the magnitude of the order parameter instantly becomes zero and the phase changes by  $2\pi$

The minimum length the superconductivity can be destroyed is the coherence length  $\xi(T)$

The minimum energy corresponds to destruction of superconductivity in a volume  $\xi(T) \sigma$ :  
 $\Delta F = B_c^2 \xi(T) \sigma$ , where  $B_c(T)$  is the critical field

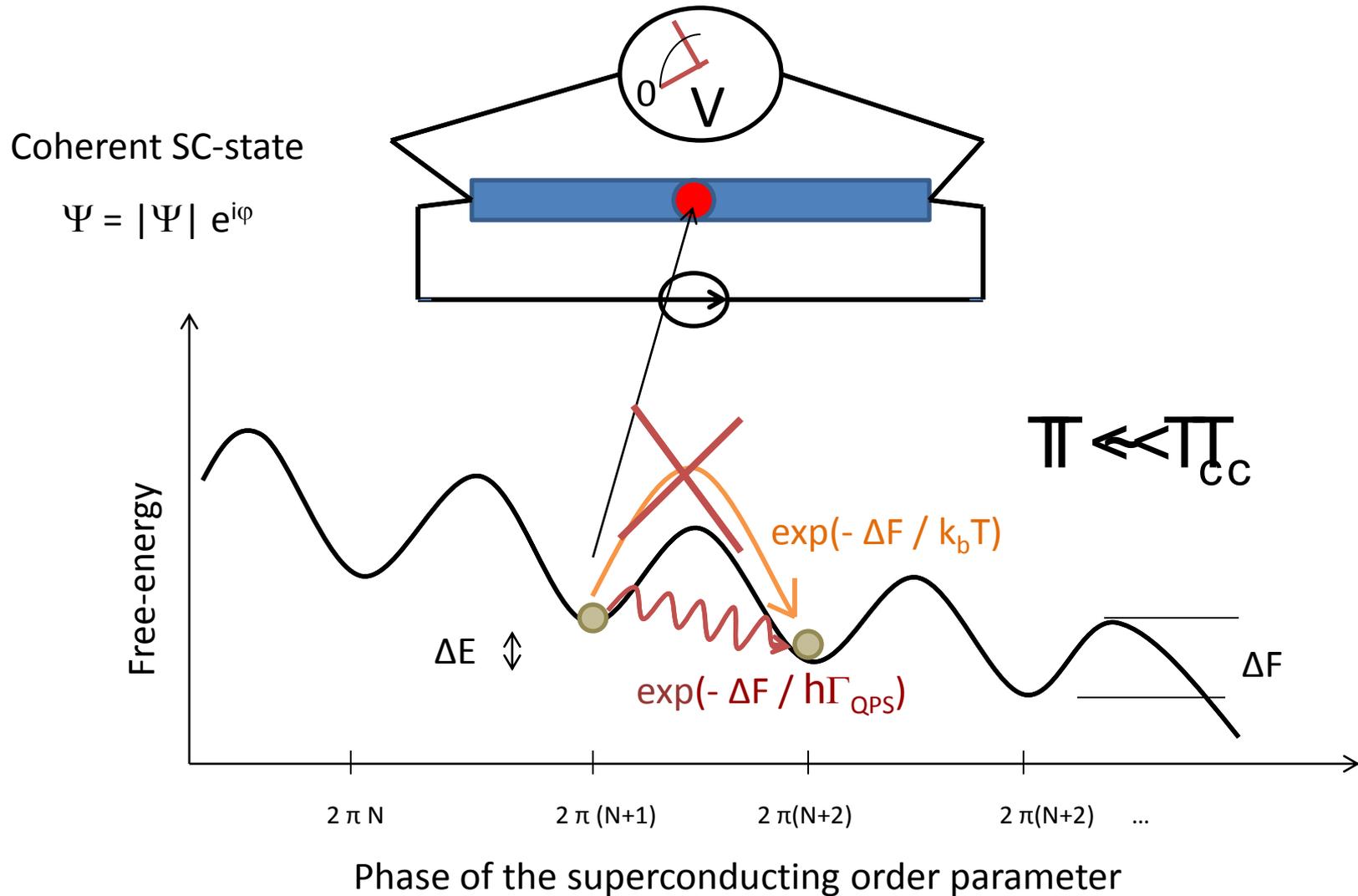
In the limit rare events the probability of the process  $P(T) \sim \exp(-\Delta F / \mathcal{E})$

**Thermal activation:  $\mathcal{E} \sim k_B T$ .**  
**Important at  $T \rightarrow T_c$ .**

**Quantum:  $\mathcal{E} \sim \Delta$ .**  
**Weak temperature dependence,**  
**exist even at  $T \rightarrow 0$**

In current state the particular manifestation of a quantum fluctuation when magnitude of the order parameter momentarily nulls and phase changes by  $\pm 2\pi$  is often called “phase slip”

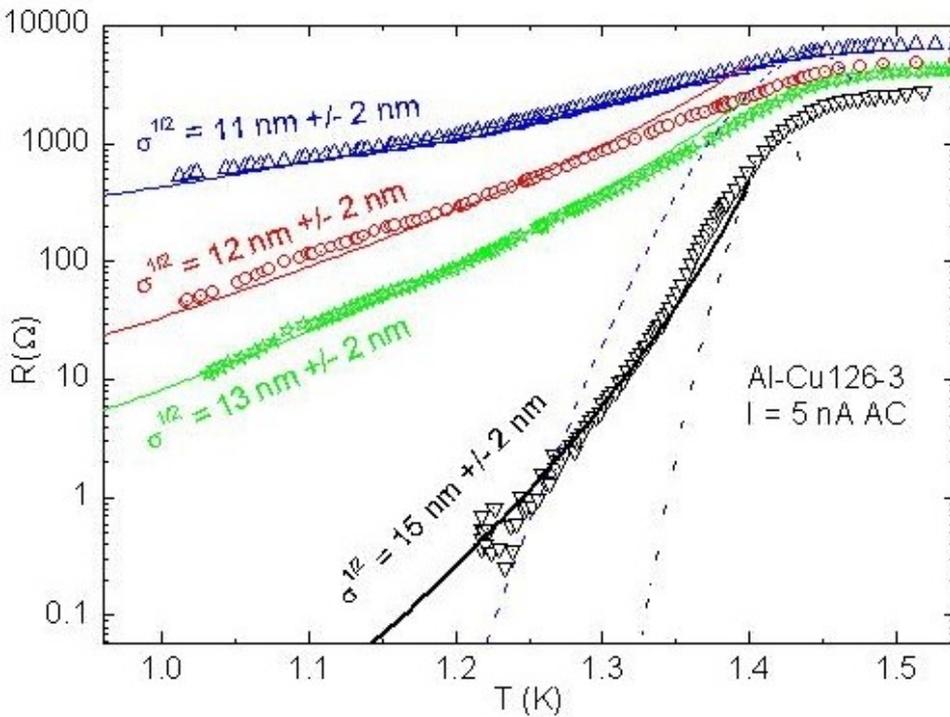
# Phase slip concept



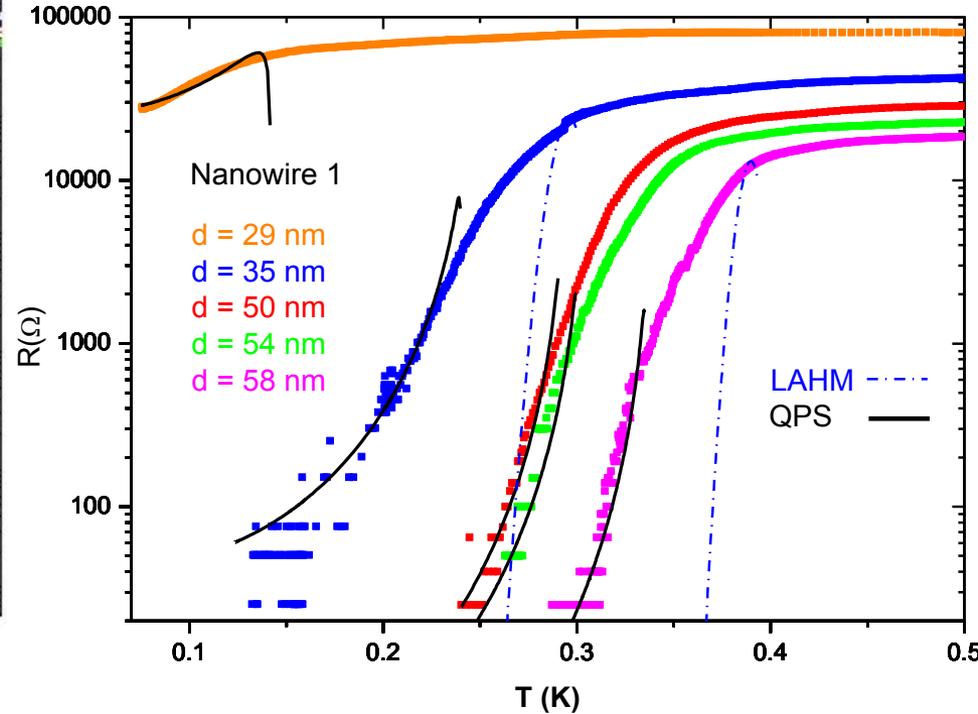
The initial and final states have different energy  $\rightarrow$  dissipation in a superconductor

# Manifestation of QPS: broadening of R(T) transition

Aluminium



Titanium



**Dashed lines: thermally activated fluctuations (LAHM) at  $T \sim T_c$**

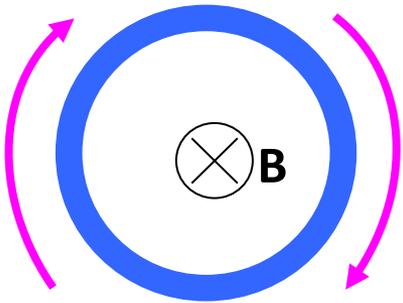
**Solid lines: quantum fluctuations (Golubev – Zaikin) model at  $T \ll T_c$**

M. Zgirski, K.-P. Riikonen, V. Touboltsev, and K. Arutyunov. // *Nano Letters* **5**, 1029--1033 (2005).

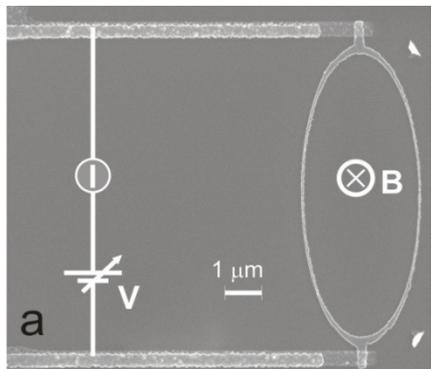
M. Zgirski, K.-P. Riikonen, V. Touboltsev and K. Yu. Arutyunov.// *PRB* **77**, 054508-1 -- 054508-6 (2008).

J. S. Lehtinen, T. Sajavara, K. Yu. Arutyunov, M. Yu. Presnjakov, and A. S. Vasiliev // *PRB* **85**, 094508 (2012).

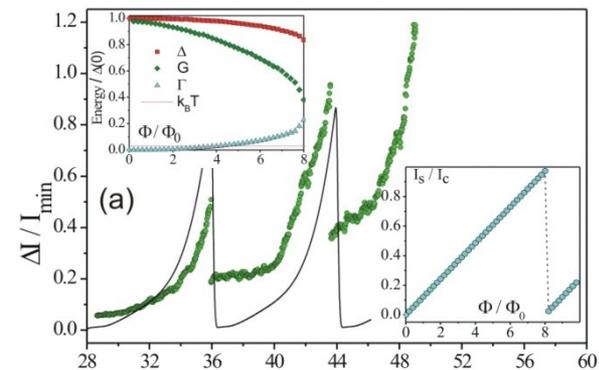
# Persistent currents in nanorings



# Amplitude and period of oscillations as function of the loop linewidth (aluminium)

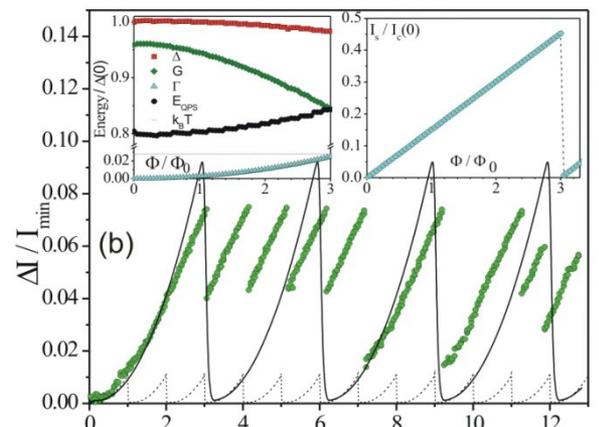


**Wide wire: classic limit**  
 110 nm x 75 nm wire,  
 $T_{bath}=65\pm 5$  mK,  $\sigma_{fit}^{1/2}=90.8$  nm



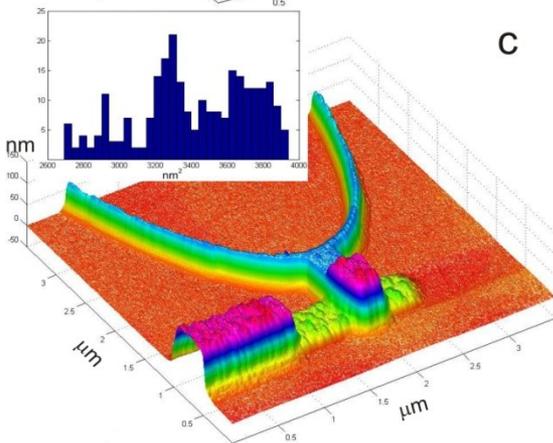
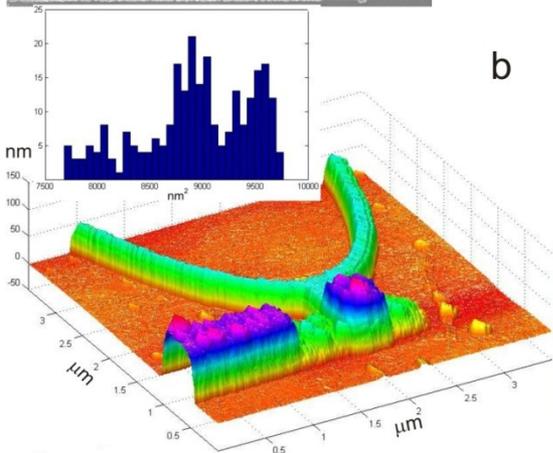
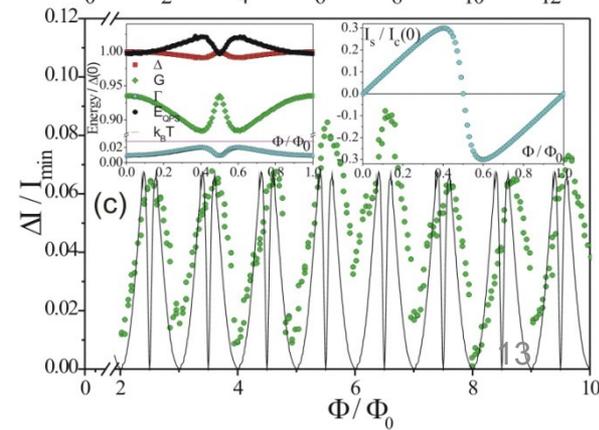
b

**Intermediate regime:**  
 $T_{bath}=52\pm 5$  mK. Solid line:  
 $\sigma_{fit}^{1/2}=12.49$  nm and  $\Delta\Phi/\Phi_0=3$   
 period, dashed line - calculated  
 $\Delta\Phi/\Phi_0=1$  oscillations in a slightly  
 narrower loop  $\sigma_{fit}^{1/2}=12.37$  nm



c

**QPS limit:**  
 same sample as in (b) further  
 gently sputtered at  $T_{bath}=54\pm 5$  mK  
 Solid line - calculations in the QPS  
 limit with  $\sigma_{fit}^{1/2}=12.15$  nm



**Fluctuations of the  
superconducting gap  
amplitude**

# Fluctuations of the order parameter

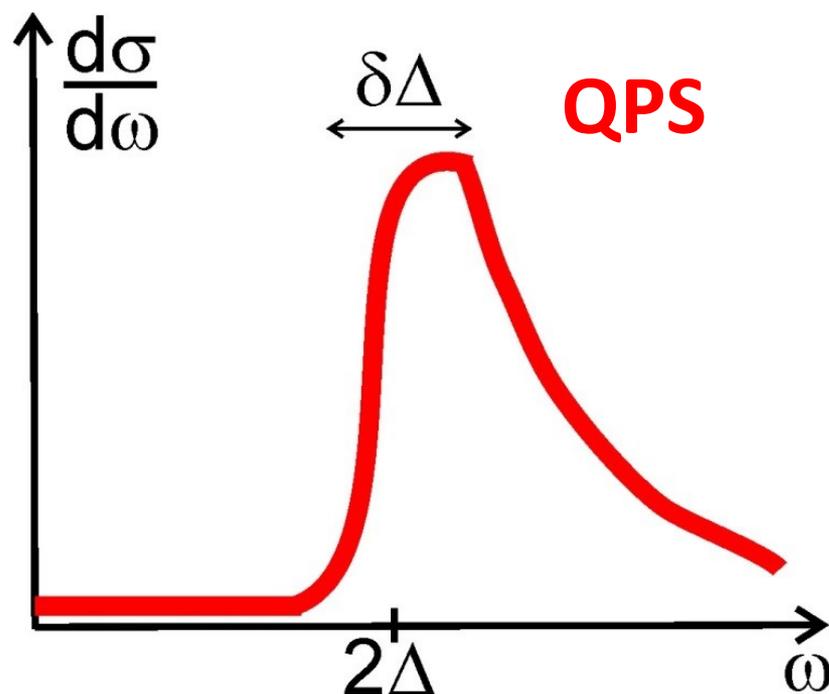
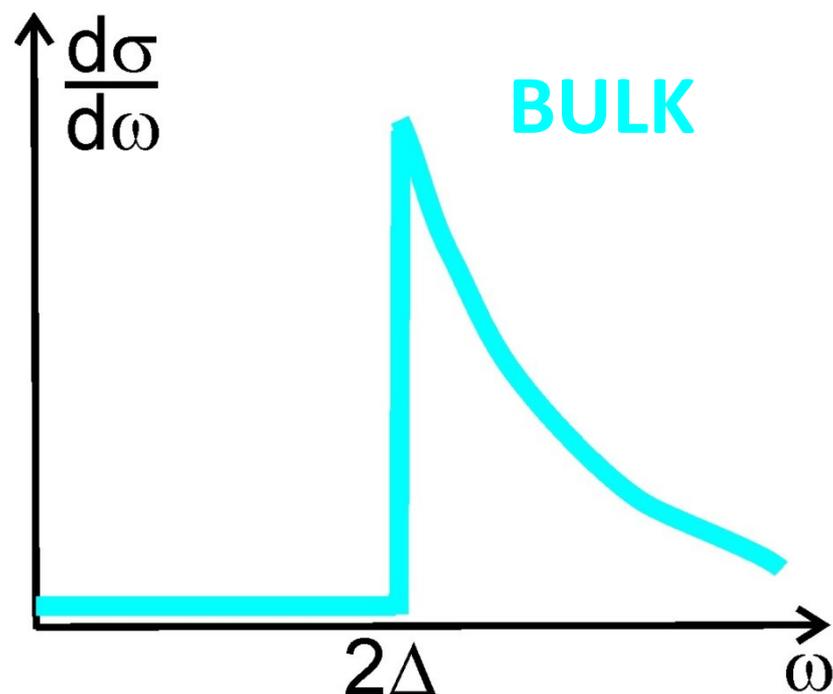
$$\Delta = \Delta_0 \exp(i\varphi)$$



# Direct determination of the fluctuating energy gap

$$\delta|\Delta| / |\Delta| \sim (S_{\text{QPS}})^{-1}$$

Superconducting gap can be associated with E/M radiation absorption threshold

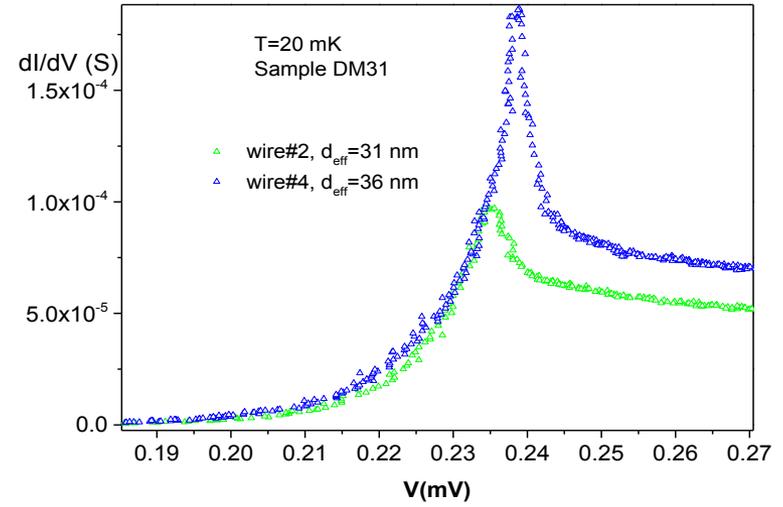
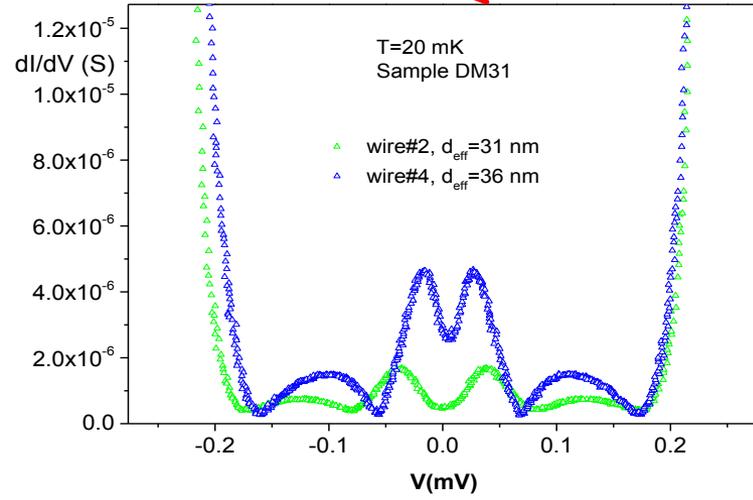
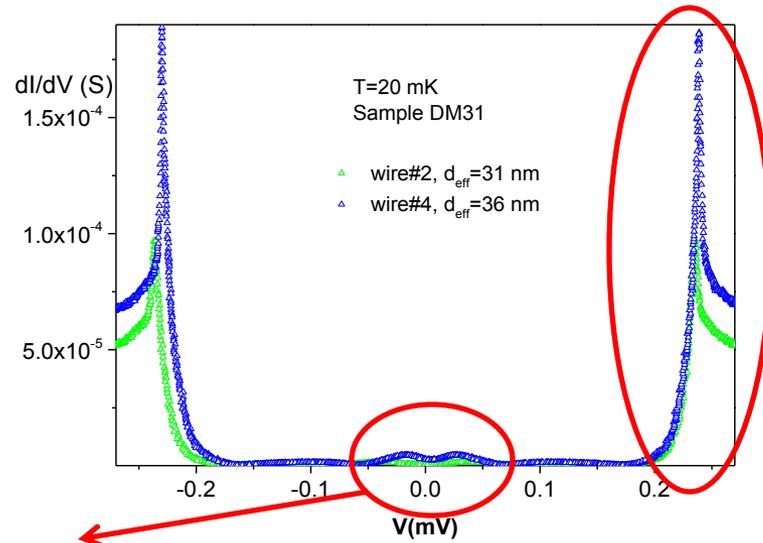


Characteristic scale

$\Delta(\text{Al}) \sim 80$  to  $100$  GHz,  $\Delta(\text{Ti}) \sim 15$  to  $25$  GHz.

$\delta\Delta / \Delta$  can reach 30% in sufficiently thin nanowires

# SIS I(V) at various diameters of nanowires



$\Delta_1 - \Delta_2$  features and

$\Delta_1 + \Delta_2$  features

Coulomb blockade and/or Josephson current (?)

**The smaller the nanowire diameter (1) the smaller the average value  $\langle \Delta \rangle$  and (2) the larger the gap smearing  $\delta \Delta$ .**

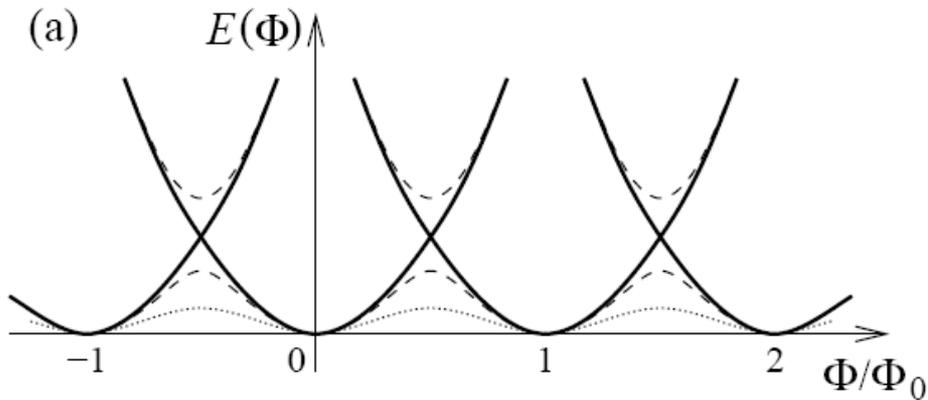
# Conclusions on physics

**Quantum fluctuations in narrow superconducting nanowires suppress 'basic' superconducting attributes manifesting as:**

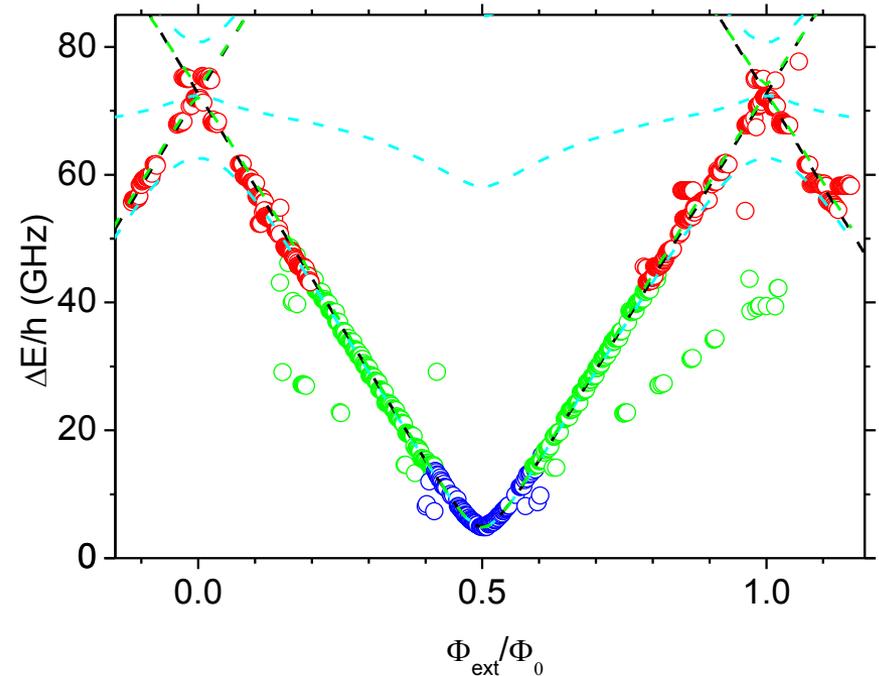
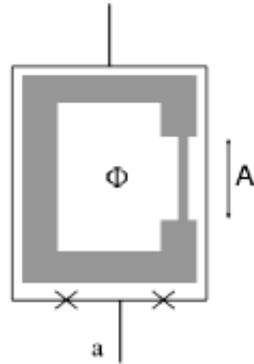
- **Broadening of the  $R(T)$  transition. In thinnest samples zero resistance is not reached even at  $T \rightarrow 0$ .**
- **Suppression of persistent currents in nanorings.**
- **Smearing of the superconducting gap edge.**

**Applications**

# Qubit



**Superconducting  
horse-shoe, shortened  
by a QPS nanowire**

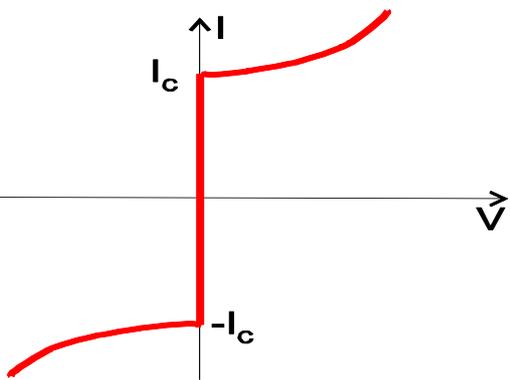


*Spectroscopy of the system  
across a wide range of flux  
and frequency.*

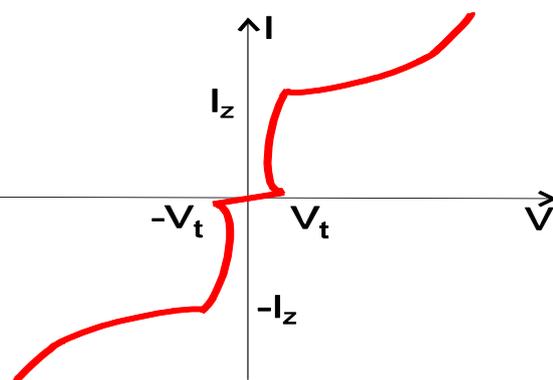
J. E. Mooij and C. J. P. M. Harmans  
**New J. Phys.** 7,219 (2005)

O. V. Astafiev, L. B. Ioffe, S. Kafanov, Yu. A. Pashkin,  
K. Yu. Arutyunov, D. Shahar, O. Cohen, and J. S. Tsai.  
**Nature** 484, 355 (2012).

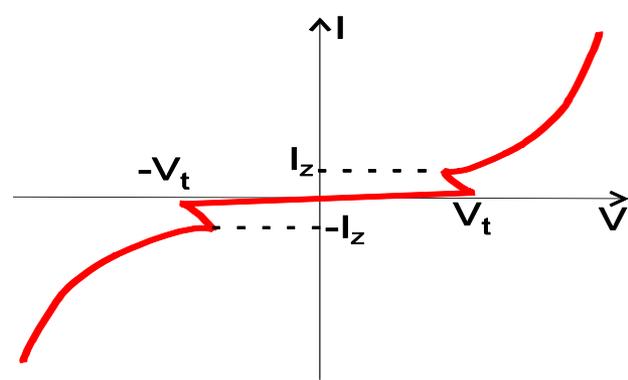
# Coulomb Effects



Thick (conventional)  
superconducting wire



Thin wire and low- $\Omega$   
environment



Very thin wire and high- $\Omega$   
environment:  
Coulomb blockade

# Quantum duality between JJ and QPS junctions

**Hamiltonian of a superconducting nanowire in the regime of quantum fluctuations:**

$$\hat{H} = \frac{E_L}{(2\pi)^2} \hat{\phi}^2 - E_{QPS} \cos(2\pi\hat{q}) + \hat{H}_{coup} + \hat{H}_{env}$$

**is dual to the corresponding Hamiltonian of a Josephson junction:**

$$\hat{H} = E_C \hat{q}^2 - E_J \cos(\hat{\varphi}) + \hat{H}_{coup} + \hat{H}_{env}$$

**with the accuracy of substitution:**  $E_C \leftrightarrow E_L, E_J \leftrightarrow E_{QPS}, \varphi \leftrightarrow \pi q / 2e$

$E_L, E_C, E_J$  and  $E_{QPS}$  are the inductive, charging, Josephson coupling and QPS energies,  $\varphi$  is phase and  $q$  is quasicharge.

**The extensively developed physics for Josephson systems can be ‘mapped’ on the superconducting nanowires in the regime of quantum fluctuations.**

*D. V. Averin and A. A. Odintsov, Phys. Lett. A* **140** (1989) 251

*J. E. Mooij and Yu. V. Nazarov, Nature Physics* **2** (2006) 169

*A. M. Hriscu and Yu. V. Nazarov, Phys. Rev. B* **83**, 174511 (2011)

**Junctionless  
transistor**

# Cooper pair transistor without a tunnel junction

*A.M. Hriscu and Yu. V. Nazarov. Phys. Rev. B 83, 174511 (2011)*

Let us consider the simplest case of a QPS Cooper pair box:

$$\hat{H}_{QPSbox} = E_c \left( \hat{Q} - \frac{q}{2e} \right)^2 + E_L \hat{\phi}^2 - E_{QPS} \cos(2\pi \hat{Q})$$

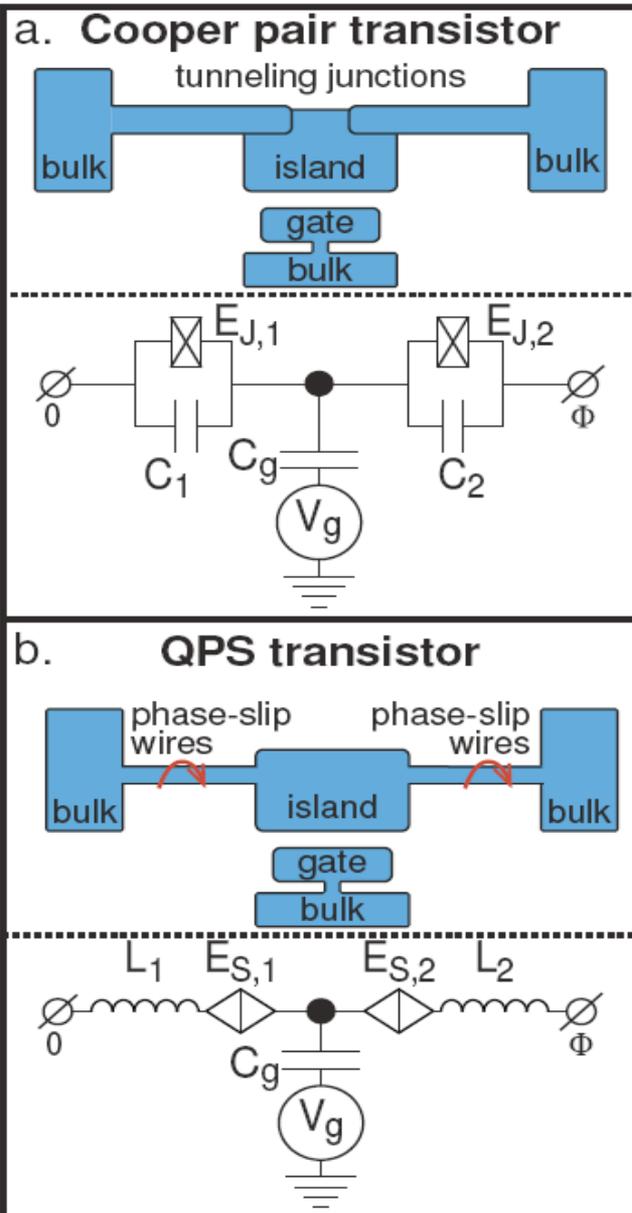
First two terms correspond to linear LC oscillator. If to shift the charge variable by  $q/2e$ , the induced charge disappears from the oscillator part, while the QPS amplitudes acquire the phase factors:

$$\hat{H}_{QPS} \psi(\phi) = -E_{QPS} e^{-\frac{i\pi q}{e}} \psi(\phi + 2\pi) - E_{QPS} e^{+\frac{i\pi q}{e}} \psi(\phi - 2\pi)$$

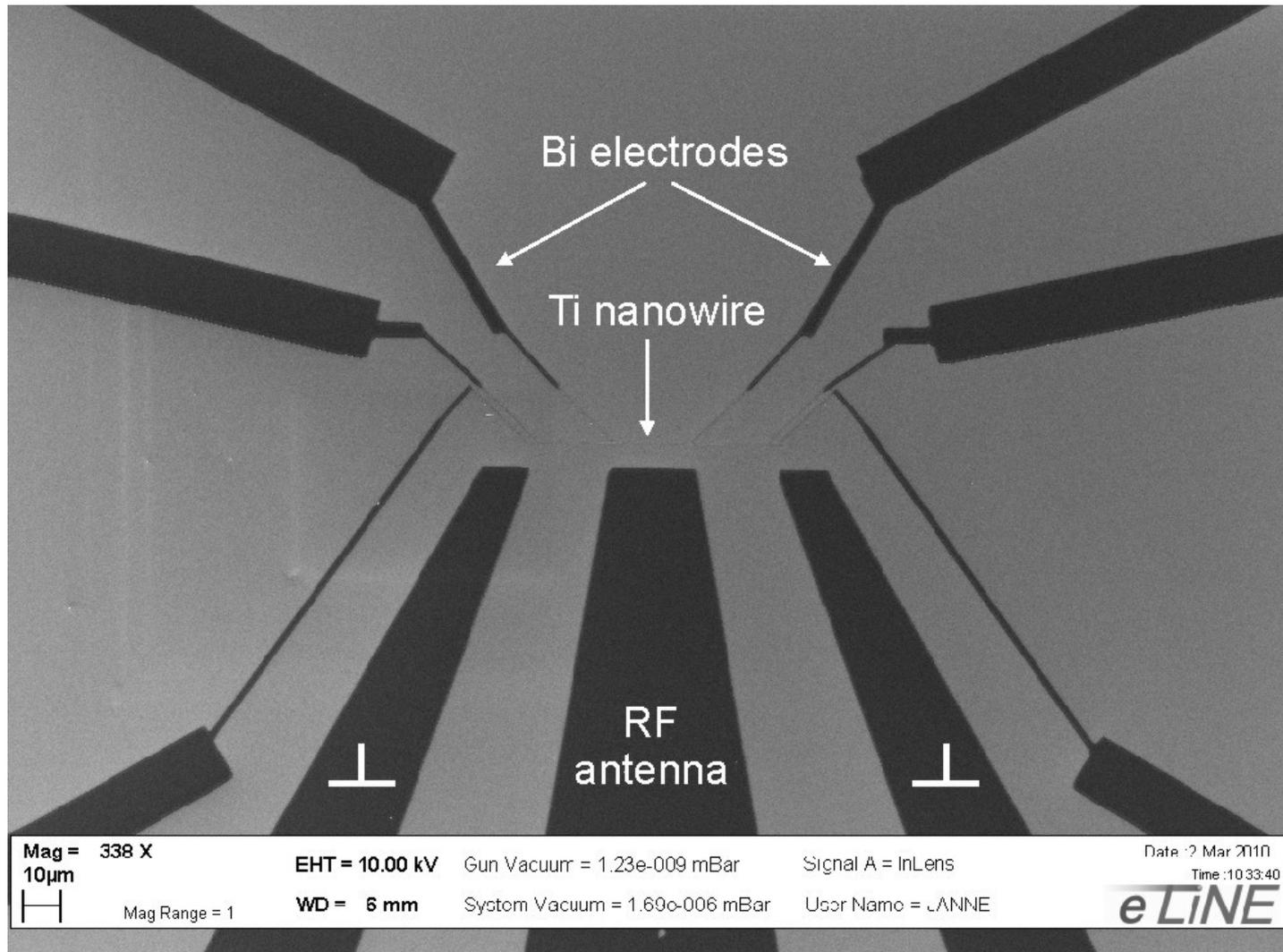
**The charge sensitivity is purely due to the coherent QPS term: the induced charge  $q/2e$  affects the interference of the phase slips with two opposite directions  $\pm 2\pi$ .**

The mandatory requirement for the charge effects observation is the high impedance of the environment: current bias  $\rightarrow$  charge is a 'good' quantum number.

**QPS is a dynamic equivalent of a conventional (static) tunnel junction.**



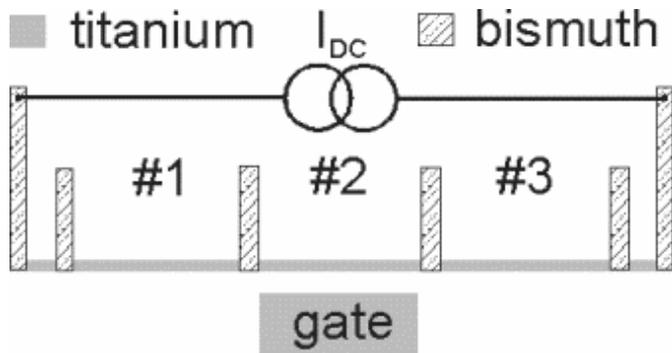
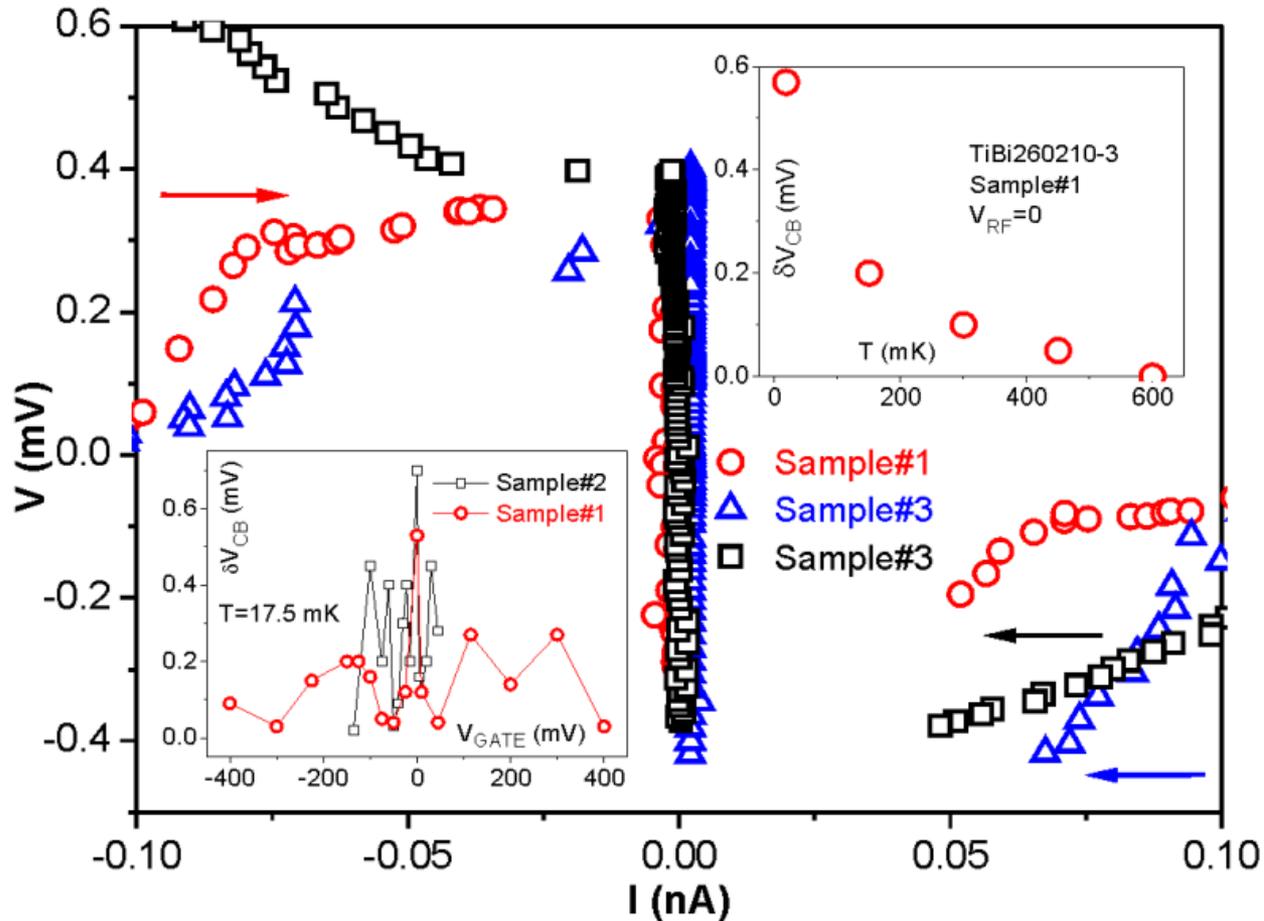
# High-Ohmic environment



Purely dissipative environment:

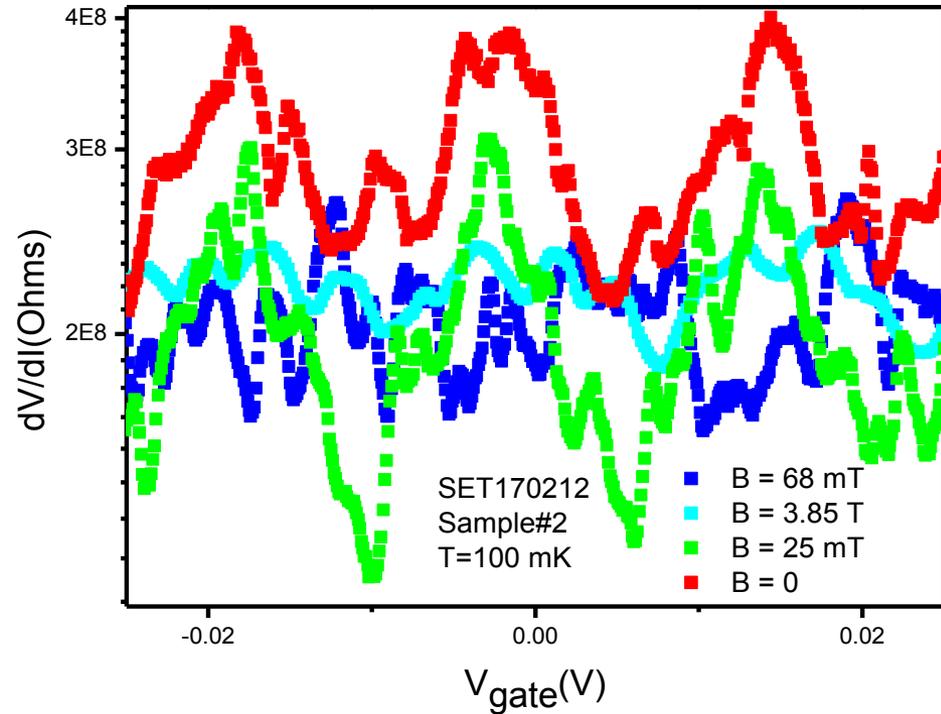
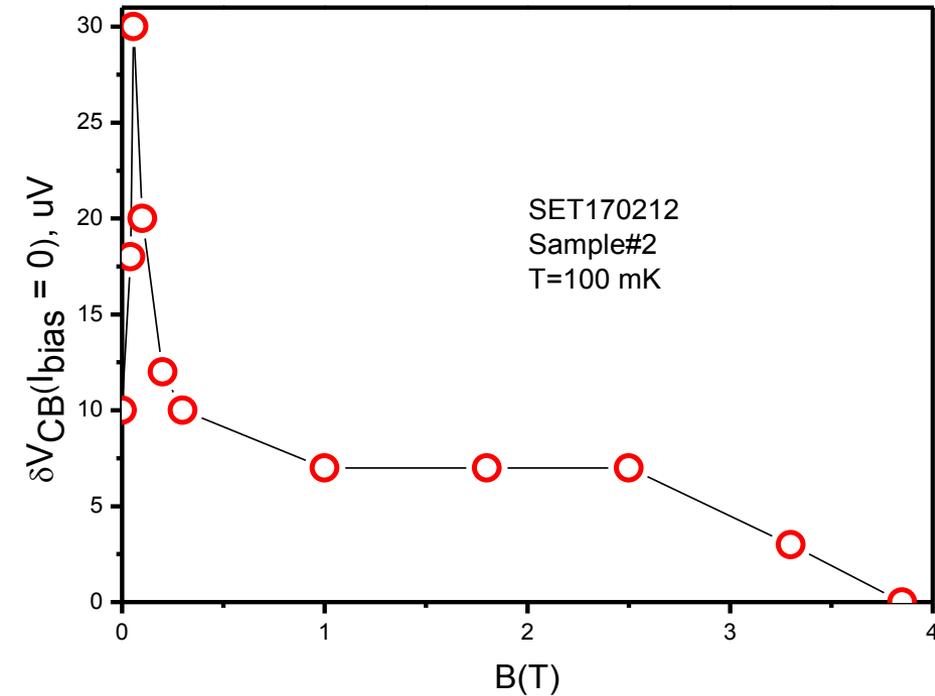
high-Ohmic normal metal probes with  $R_{\text{probe}}$  up to 10 MΩ

# 24 nm titanium nanowire, 10 M $\Omega$ contacts



All three neighboring parts of the same multiterminal structure demonstrate the same value of the Coulomb gap. The effect disappears above  $T_c$  and/or  $H_c$ .

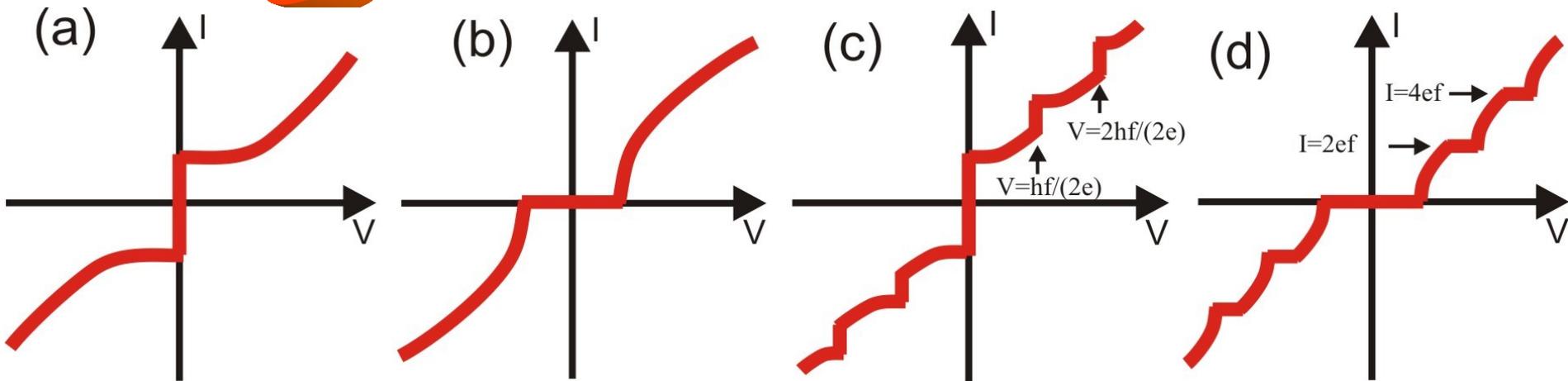
# Magnetic field dependencies



At a given (low) temperature the Coulomb gap and the gate modulation disappear above certain magnetic field (the  $B_c$  of Ti?).

Electron transport in a Cooper pair transistor is a periodic process corresponding to cyclic charging/discharging of the central island by  $2e$ . Synchronisation of the process with external drive should lead to resonance.

# Bloch Oscillations



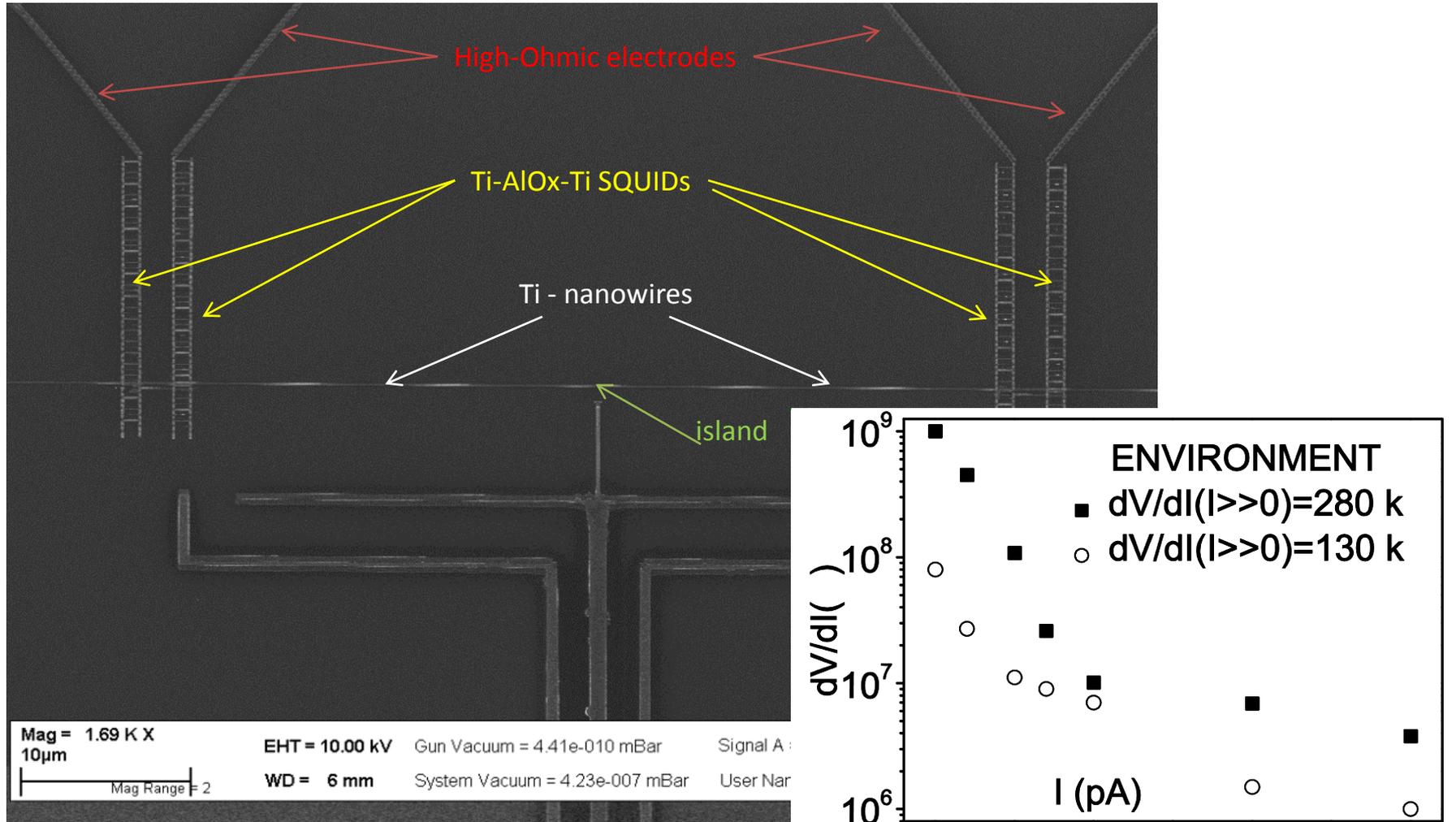
Josephson junction:  
critical current

QPS wire:  
critical voltage

Josephson junction:  
voltage steps  
(Shapiro effect)

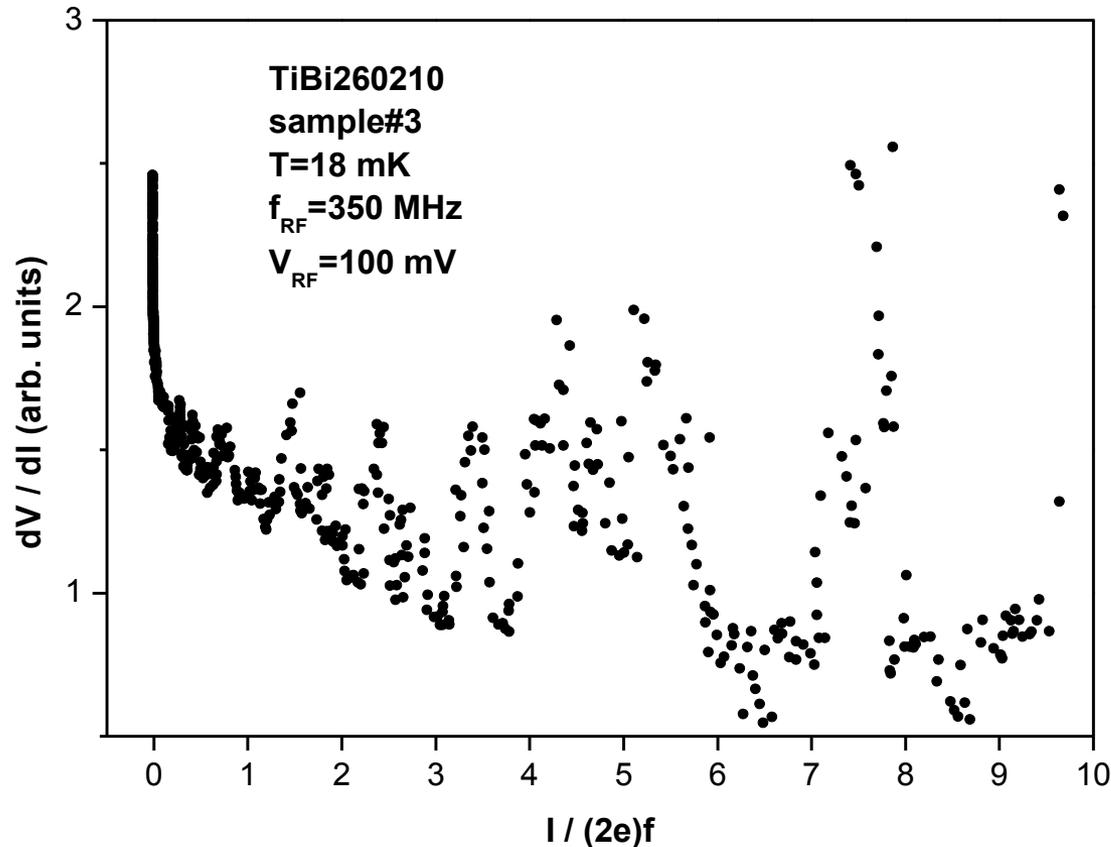
QPS wire:  
current steps

# Hybrid high impedance environment



High-Ohmic normal metal probes with  $R_{\text{probe}}$  up to 500 K $\Omega$  and 1D array of SQUIDs.

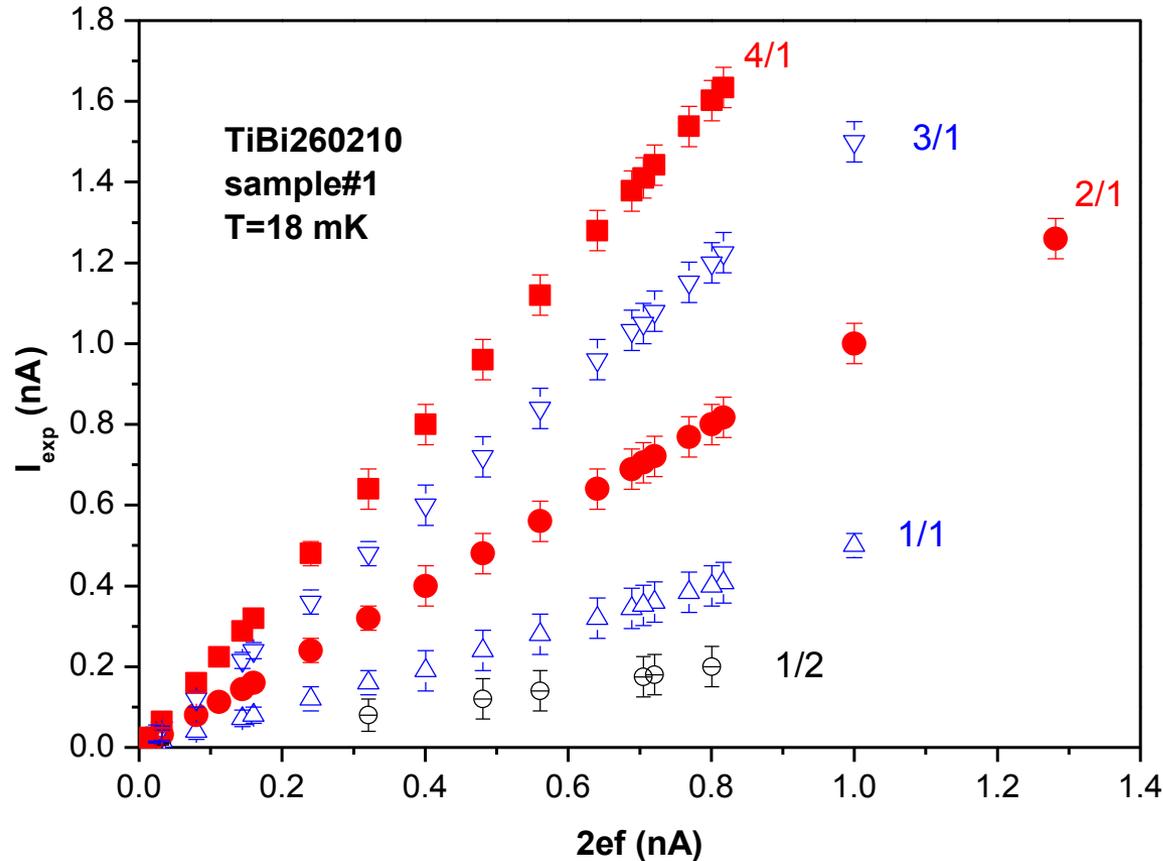
# Ti-nanowire with high-Ohmic environment



*$dV/dI$  at  $f_{RF}=350$  MHz and amplitude 100 mV. Current is normalized by  $(2e) \times f_{RF}$ . One can distinguish steps with quantum number  $n \leq 8$ .*

$$\text{Universal relation } I(n) = (2e) * n * f_{RF}$$

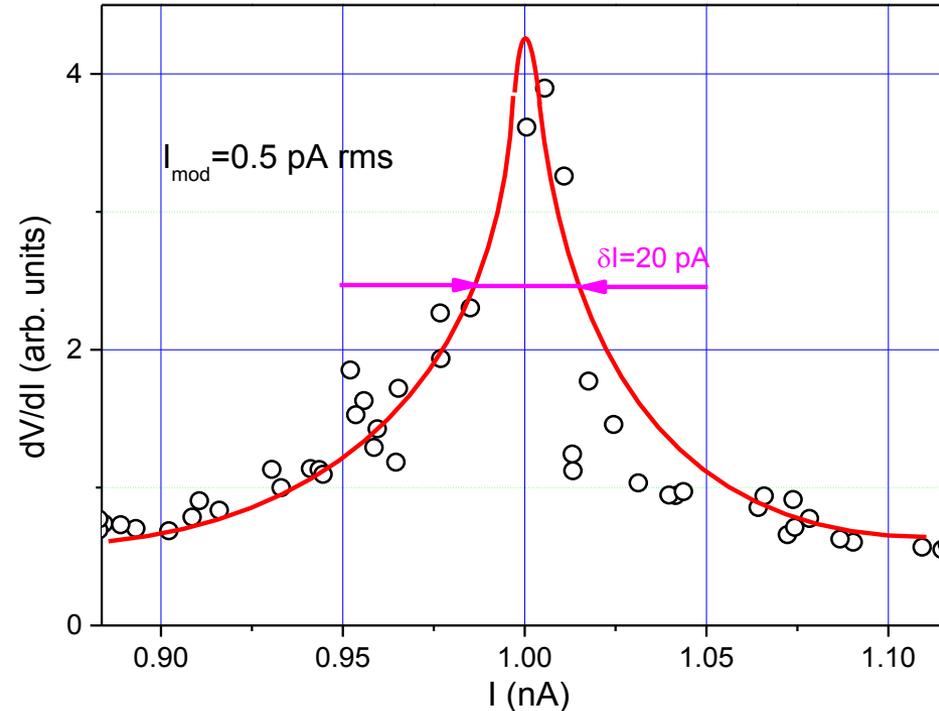
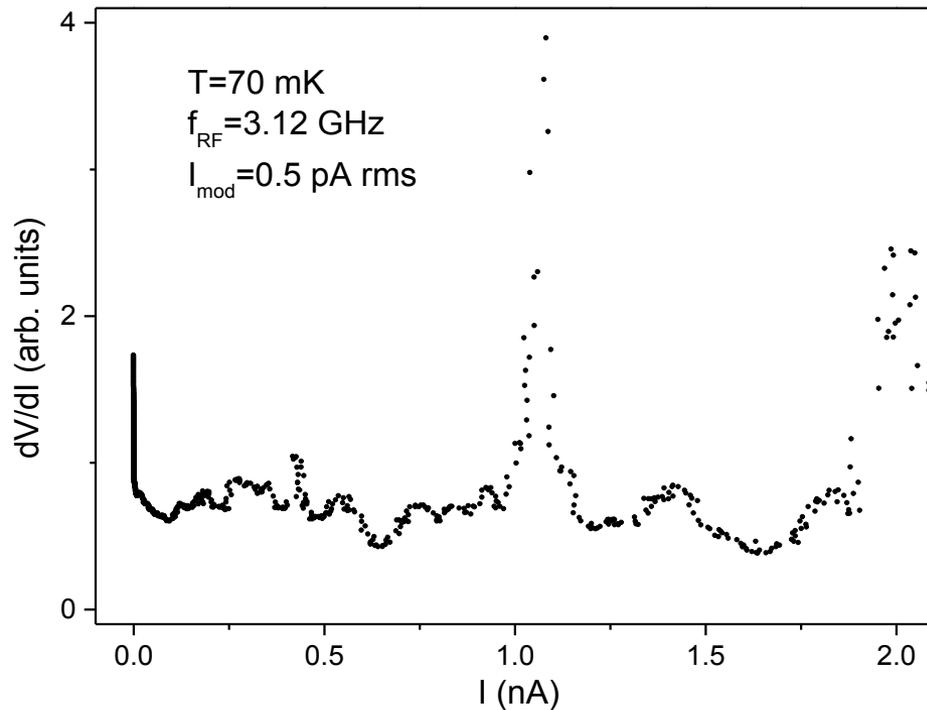
# Positions of the RF-induced steps



Red symbols correspond to Bloch steps ( $2e$ ), blue – single electron ( $e$ ), black – 1st subharmonic of single electron oscillations.

**Proof-of-principle demonstration of a quantum standard for electric current**

# Width of the Bloch step



**Demonstrated accuracy is +/- 2%.**

**Further accuracy improvement is mandatory for practical metrology.**

# Digest on Coulomb effects in QPS

- Quantum phase slip and Josephson tunneling are the phenomena described by identical Hamiltonians: the quantum dynamics is indistinguishable.
- Superconducting nanowire (homogeneous!), in the regime of QPS, is the dynamic equivalent of a Josephson junction.
- Containing no static (in space and time) junctions, QPS nanowire can sustain much higher currents and has no undesired two-level 'fluctuators' present in tunnel contacts.
- All phenomena, observable in Josephson systems, can be observed in QPS nanowires.

**Thank you!**

# Summer University in Moscow

## Quantum Technologies

"Quantum Technologies" program provides the unique opportunity to get introduction to the quickly developing interdisciplinary topics of quantum solid state physics, quantum photonics and nanotechnology. The program includes both the introductory and the rather advanced courses at the very frontier of modern science. The invited lecturers are the world-class leading experts in corresponding fields. The training is addressed to students of bachelor / magister level who have affinity with physics, electronics, IT and computer sciences. The courses might appear attractive also to researchers in natural sciences such as chemistry and biology who wish to advance their knowledge in various phenomena related to nanoscience.

The program at the Summer University is the pilot project announcing the new English-speaking full-scale MSc program "Quantum Information Technologies (iQT)" to start at Higher School of Economics in 2017. The objective of that program is to train specialists of the highest qualification in the field of transmission, storage and processing of information using the most advanced methods of quantum physics and quantum communication. The announced courses of Summer University-2016 have been prepared by the core lecturers of the upcoming iQT program, and represent the selected and shortened versions of MSc training in quantum communication.

<https://www.hse.ru/international/summer/quantum>